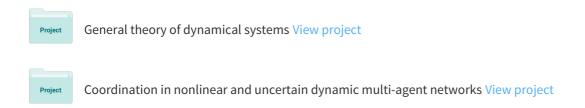
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Model reduction of a flexible-joint robot: a port-Hamiltonian approach

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Abstract: In this paper we explore the methodology of model order reduction based on singular perturbations for a flexible-joint robot within the port-Hamiltonian framework. We show that a flexible-joint robot has a port-Hamiltonian representation which is also a singularly perturbed ordinary differential equation. Moreover, the associated reduced slow subsystem corresponds to a port-Hamiltonian model of a rigid-joint robot. To exploit the usefulness of the reduced models, we provide a numerical example where an existing controller for a rigid robot is implemented.

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1. INTRODUCTION

This document explores the model order reduction of a port-Hamiltonian (PH) system based on singular perturbations. The case study is a flexible-joint robot. By a series of transformations we show that a model for a flexible-joint robot can be written as 1) a PH system and 2) a singularly perturbed ordinary differential equation. Afterwards we show the effect of a composite control based on the reduced subsystems.

In the Euler-Lagrange (EL) framework, position control robotic systems have been thoroughly discussed in e.g., Canudas-de Wit et al. (1996); Murray et al. (1994); Ortega et al. (1998); Spong et al. (2006). In such a framework, the control design is based on selecting a suitable storage function that ensures position control. However, the desired storage function under the EL framework does not qualify as an energy function in any physical meaningful sense as stated in Canudas-de Wit et al. (1996); Ortega et al. (1998).

The PH modeling framework of van der Schaft and Maschke (2003); van der Schaft (2000) has received a considerable amount of interest in the last decade because of its insightful physical structure. It is well known that a large class of (nonlinear) physical systems can be described in the PH framework. The popularity of PH systems can be largely accredited to its application to analysis and control design of physical systems, e.g. Duindam et al. (2009); Fujimoto and Sugie (2001); van der Schaft and Maschke (2003); van der Schaft (2000). Control laws in the PH framework are derived with a clear physical interpretation via direct shaping of the closed-loop energy, interconnection, and dissipation structure, see Duindam et al. (2009); van der Schaft (2000).

On the other hand, model order reduction plays a crucial role in control design as well. Being able to synthesize controllers with a low number of variables is always more convenient. Moreover, it is convenient as well to use reduction methods that preserve the structure of the original system, van der Schaft and Polyuga (2009); Polyuga and van der Schaft (2010); Scherpen and van der Schaft (2008). One of the several model reduction methods is based on singular perturbations, which is often applied to systems with two or more time-scales. It is known that under certain hyperbolicity properties (see Section 2) it is possible to obtain reduced models corresponding to systems in distinct time scales. The behavior of the full system can be inferred by an analysis of the reduced systems.

In the following sections we present a case study where model order reduction, based on singular perturbation, is applied to a PH system. In Section 2, we present a general background in the PH framework, especially for a class of standard mechanical systems. Furthermore, we recall the results of Viola et al. (2007) to equivalently describe the original PH system in a PH form which has a constant mass-inertia matrix in the Hamiltonian via a change of variables. This will prove helpful when writing the model of a flexible-joint robot. In Section 2 we also recall some basic properties of slow-fast systems. Afterwards in Section 3 we present a model of a flexible-joint robot that has a PH and a slow-fast structure. The corresponding reduced subsystem is the model of a rigid robot. We conclude this document with a simulation of a 2R flexible-joint robot for which a controller is designed based on the reduced models.

2. PRELIMINARIES

In this section we present the PH formalism for a class of standard mechanical systems. Additionally, we recapitulate the results of Fujimoto and Sugie (2001) in terms of generalized coordinates transformations for PH systems. We also recall the results of Viola et al. (2007) to transform the original system into PH form with a constant mass-

inertia matrix. Next, we include a brief introduction on slow-fast systems.

2.1 Port-Hamiltonian Systems

The PH framework is based on the description of systems in terms of energy variables, their interconnection structure, and power ports. PH systems include a large family of physical nonlinear systems. The transfer of energy between the physical system and the environment is given through energy elements, dissipation elements and power preserving ports, see Duindam et al. (2009); van der Schaft and Maschke (2003); van der Schaft (2000).

A class of PH system, introduced by van der Schaft and Maschke (2003), is described by

$$\Sigma = \begin{cases} \dot{x} = \left[J(x) - R(x)\right] \frac{\partial H(x)}{\partial x} + g(x)w \\ y = g(x)^{\top} \frac{\partial H(x)}{\partial x} \end{cases}$$

with states $x \in \mathbb{R}^{\mathcal{N}}$, skew-symmetric interconnection matrix $J(x) \in \mathbb{R}^{\mathcal{N} \times \mathcal{N}}$, positive semi-definite damping matrix $R(x) \in \mathbb{R}^{\mathcal{N} \times \mathcal{N}}$, and Hamiltonian $H(x) \in \mathbb{R}$. The matrix $g(x) \in \mathbb{R}^{\mathcal{N} \times \mathcal{M}}$ weights the action of the control inputs $w \in \mathbb{R}^{\mathcal{M}}$ on the system, and $w, y \in \mathbb{R}^{\mathcal{M}}$ with $\mathcal{M} \leq \mathcal{N}$, form a power port pair.

In this preliminary, we restrict the analysis to the a class of standard mechanical systems with n ($\mathcal{N}=2n$) degrees of freedom (dof),

$$\begin{bmatrix} \dot{q} \\ \dot{p} \end{bmatrix} = \begin{bmatrix} 0_{n \times n} & I_{n \times n} \\ -I_{n \times n} & -D(q, p) \end{bmatrix} \begin{bmatrix} \frac{\partial H(q, p)}{\partial q} \\ \frac{\partial H(q, p)}{\partial p} \end{bmatrix} + \begin{bmatrix} 0_{n \times n} \\ G(q) \end{bmatrix} u$$

$$y = G(q)^{\top} \frac{\partial H(q, p)}{\partial p}$$
(1)

with generalized configuration coordinates $q \in \mathbb{R}^n$, generalized momenta $p \in \mathbb{R}^n$, damping matrix $D\left(q,p\right) \in \mathbb{R}^{n \times n}$, where $D\left(q,p\right) = D\left(q,p\right)^{\top} \geq 0$, output $y \in \mathbb{R}^n$, input $u \in \mathbb{R}^n$, and the input matrix $G\left(q\right) \in \mathbb{R}^{n \times n}$. The Hamiltonian function of (1) is given by

$$H(q,p) = \frac{1}{2} p^{\top} M^{-1}(q) p + V(q)$$
 (2)

where $M\left(q\right)=M^{\top}\left(q\right)>0$ is the $n\times n$ inertia (generalized mass) matrix, and $V\left(q\right)$ is the potential energy.

2.2 Nonconstant to constant mass-inertia matrix

Consider the class of standard mechanical systems in (1) with a nonconstant mass-inertia matrix M(q). The aim of this section is to transform (1) into a PH system with a constant mass-inertia matrix by a generalized canonical transformation, see Fujimoto and Sugie (2001); Viola et al. (2007).

Consider the system in (1) with nonconstant M(q), and a coordinate transformation $(\bar{q}, \bar{p}) = \Phi(q, p)$ where

$$\Phi\left(q,p\right) = \begin{pmatrix} q - q^* \\ T\left(q\right)^{-1} p \end{pmatrix} = \begin{pmatrix} q - q_d \\ T\left(q\right)^{\top} \dot{q} \end{pmatrix} = \begin{pmatrix} \bar{q} \\ \bar{p} \end{pmatrix} \quad (3)$$

with $q^* \in \mathbb{R}^n$ being a constant position vector, and T(q) a lower triangular matrix such that

$$T(q) = T\left(\Phi^{-1}(\bar{q}, \bar{p})\right) = \bar{T}(\bar{q})$$

and

$$M\left(q\right) = T\left(q\right)T\left(q\right)^{\top} = \bar{T}\left(\bar{q}\right)\bar{T}\left(\bar{q}\right)^{\top}$$

Consider now the Hamiltonian H(q,p) as in (2). Using (3), the new function $\bar{H}(\bar{x}) = H(\Phi^{-1}(\bar{x}))$ and $\bar{V}(\bar{q}) = V(\Phi^{-1}(\bar{q}))$ read as

$$\bar{H}\left(\bar{x}\right) = \bar{H}\left(\bar{q}, \bar{p}\right) = \frac{1}{2} \bar{p}^{\top} \bar{p} + \bar{V}\left(\bar{q}\right)$$

Using this Hamiltonian and the coordinate transformation in (3), our system under consideration in (1) can be rewritten, as in van der Schaft and Maschke (2003), as follows

$$\begin{bmatrix} \dot{\bar{q}} \\ \dot{\bar{p}} \end{bmatrix} = \begin{bmatrix} 0_{n \times n} & \bar{T}^{-\top} \\ -\bar{T}^{-1} & \bar{J}_2 - \bar{D} \end{bmatrix} \begin{bmatrix} \frac{\partial \bar{H} (\bar{q}, \bar{p})}{\partial \bar{q}} \\ \frac{\partial \bar{H} (\bar{q}, \bar{p})}{\partial \bar{p}} \end{bmatrix} + \begin{bmatrix} 0_{n \times n} \\ \bar{G} \end{bmatrix} \bar{v} (4)$$
$$\bar{y} = \bar{G}^{\top} \frac{\partial \bar{H} (\bar{q}, \bar{p})}{\partial \bar{p}}$$

with a new input $\bar{v} \in \mathbb{R}^n$, and where the skew-symmetric matrix $\bar{J}_2 = \bar{J}_2(\bar{q}, \bar{p})$ takes the form

$$\bar{J}_{2}(\bar{q},\bar{p}) = \frac{\partial \left(\bar{T}^{-1}\bar{p}\right)}{\partial \bar{q}}\bar{T}^{-\top} - \bar{T}^{-1}\frac{\partial \left(\bar{T}^{-1}\bar{p}\right)}{\partial \bar{q}}^{\top} \tag{5}$$

with

$$(q,p) = \Phi^{-1}(\bar{q},\bar{p})$$

together with the matrix $\bar{D}\left(\bar{q},\bar{p}\right)$, and the input matrix $\bar{G}\left(\bar{q}\right)$ given by

$$\bar{D}\left(\bar{q},\bar{p}\right) = \bar{T}\left(\bar{q}\right)^{-1} D\left(\Phi^{-1}\left(\bar{q},\bar{p}\right)\right) \bar{T}\left(\bar{q}\right)^{-\top}$$
$$\bar{G}\left(\bar{q}\right) = \bar{T}\left(\bar{q}\right)^{-1} G\left(\bar{q}\right)$$

In (4) and (5) we have left out the argument \bar{q} of \bar{T} (\bar{q}) for notational simplicity.

Remark 1. The change of coordinates described in this section is fundamental for the result of Section 3. Through this change of coordinates it is possible to write the model of a flexible joint robot as a PHS and a slow-fast system. Up to the authors' experience, if such a transformation is not performed, a slow-fast PHS is much more difficult to obtain.

2.3 Slow-fast systems

By a slow-fast system (SFS) we mean a singularly perturbed ordinary differential equation of the form

$$\dot{x} = f(x, z, \varepsilon)
\varepsilon \dot{z} = q(x, z, \varepsilon),$$
(6)

where $x \in \mathbb{R}^m$, $z \in \mathbb{R}^n$ and $\varepsilon > 0$ is a small parameter, i.e., $\varepsilon \ll 1$, and where f and g are smooth functions. Note that due to the presence of the parameter ε , the variable z evolves much faster than x. Then, x and z are called the slow and the fast variables respectively. For $\varepsilon \neq 0$ the new time parameter $\tau = t/\varepsilon$ is defined. As a consequence, (6) is rewritten as

$$x' = \varepsilon f(x, z, \varepsilon)$$

$$z' = g(x, z, \varepsilon).$$
(7)

where the prime denotes the derivative with respect to the re-scaled time τ . Note that, for $\varepsilon > 0$ and g not identically zero, the systems (6) and (7) are equivalent. A first approach to study the dynamics of slow-fast systems is to analyze the limit $\varepsilon \to 0$ of (6) and (7). These limits

correspond to a Differential Algebraic Equation (DAE)

$$\dot{x} = f(x, z, 0)
0 = q(x, z, 0),$$
(8)

and to the layer equation

$$x' = 0$$

 $z' = q(x, z, 0).$ (9)

Associated to those two limit equations, the critical manifold is defined as follows.

Definition 1. The critical manifold S is defined as

$$S = \{(x, z) \in \mathbb{R}^m \times \mathbb{R}^n \mid g(x, z, 0) = 0\}.$$

Note that the critical manifold S serves as the phase-space of the DAE (8) and as the set of equilibrium points of the layer equation (9).

If S is a set of hyperbolic points of (9), then S is called normally hyperbolic. Geometric Singular Perturbation Theory (GSPT), see e.g. Fenichel (1979); Kaper (1999); Jones (1995) shows that compact, normally hyperbolic invariant manifolds persist under C^1 -small perturbations. In the present context, this means that if $S_0 \subseteq S$ is a compact normally hyperbolic set, then, for $\varepsilon > 0$ sufficiently small, there exists a normally hyperbolic invariant manifold S_{ε} of the slow-fast system (6) which is diffeomorphic to S_0 and lies within distance of order $O(\varepsilon)$ from S_0^{-1} . This implies that the flow along S_{ε} is approximately given by the flow of the DAE (8), along S_0 .

Observe that if the matrix $\partial_z q(x,z,0)$ is regular on S, then by the implicit function theorem there exists a smooth function ϕ such that S is given as a graph $z = \phi(x)$. Therefore, the flow along the critical manifold S is defined by

$$\dot{x} = f(x, \phi(x), 0), \tag{10}$$

which is called the reduced slow vector field.

Let $x(t,x_0)$ be the flow defined by (10). The arguments of GSPT imply that the flow along the invariant manifold S_{ε} of (6) is given by $x(t,x_0) + O(\varepsilon)$. Moreover, assume that S is a set of exponentially stable equilibrium points of the layer equation (9). Then, there is a neighborhood Dof S where all trajectories with initial condition in D are exponentially attracted to the invariant manifold S_{ε} .

In the context of control systems, the hyperbolicity property has been essential in the design of controllers based on model reduction. This is mainly because if the system is does not have the hyperbolicity property, it cannot be decoupled into two (slow and fast) reduced subsystems.

Let us briefly recall the basic design methodology of composite control, for more details see Kokotovic et al. (1976); Kokotovic (1984); Kokotovic et al. (1986). Suppose we now study the control system

$$\dot{x} = f(x, z, u, \varepsilon)
\varepsilon \dot{z} = g(x, z, u, \varepsilon).$$
(11)

The strategy is to consider the reduced systems

$$\dot{x} = f(x, \phi(x), u_s, 0), \tag{12}$$

and

$$z' = g(x, z, u_s + u_f, 0), \tag{13}$$

where $u_s = u_s(x)$ denotes the controller for the reduced slow system (12) and $u_f = u_f(x,z)$ the controller for the fast subsystem (13). The idea is to design controllers u_s and u_f that: 1) make the origin of the slow subsystem (12) exponentially stable, and 2) make the critical manifold $z = \phi(x)$ exponentially stable. Then the controller u for the slow-fast system (11) is designed as $u = u_s + u_f$.

3. SLOW-FAST PORT-HAMILTONIAN MODEL

In this section, we derive a PH model of a flexible-joint robot which also has a slow-fast structure (6). In this way, the justification of designing controllers based on the reduced models is immediate. To start, let us make the following standard assumptions (Spong (1987); De Luca (2014)):

- All joints are of rotatory type
- The relative displacement (deflection) at each joint is small. Therefore we use a linear model for the springs.
- The *i*-th motor, which drives the *i*-th link, is mounted at the (i-1)-th link.
- The center of mass of the motors are located along the rotation axes.
- The angular velocity of the motors is due only to their own spinning.

We denote by $q_1 \in \mathbb{R}^n$ the links' angular positions and by $q_2 \in \mathbb{R}^n$ the motors' angular displacement.

To obtain the Hamiltonian associated to the Energies: *n*-flexible-joint robot, let us first list the involved energies:

- Link's kinetic energy: $K_l(q_1,\dot{q}_1) = \frac{1}{2}\dot{q}_1^T M_l(q_1)\dot{q}_1$, where
- $M_l(q_1)=M_l(q_1)^T>0.$ Motor's kinetic energy: $K_m(\dot{q}_2)=\frac{1}{2}\dot{q}_2^TI\dot{q}_2,$ where I= $I^T > 0$ denotes the motors' inertia.
- Potential energy due to gravity: $P_g(q_1) = \sum_{i=1}^{n} (P_{g,l_i}(q_1) + P_{g,m_i}(q_1))$, where P_{g,l_i} and P_{g,m_i} are the potential energies due to the links and the motors, respectively.
- Potential energy due to joint stiffness: $P_s(q_1, q_2) =$ $\frac{1}{2}(q_1-q_2)^TK(q_1-q_2)$, where $K \in \mathbb{R}^{n \times n}$ is a symmetric, positive definite matrix of stiffness coefficients.

Now, let us assume that the stiffness coefficient is much higher than any other parameter of the system. This is we let K be defined as $K = 1/\varepsilon \mathbb{I}_n$, where \mathbb{I}_n denotes the n-dimensional identity matrix. Next, as it is customary, e.g. Spong (1987), let us define new coordinates as $(q_1, \varepsilon z) = (q_1, q_1 - q_2)$, and denote by $q_{\varepsilon} \in \mathbb{R}^{2n}$ the generalized coordinates $q_{\varepsilon}=(q_1,z)$. Then the corresponding Hamiltonian $H_{\varepsilon}=H_{\varepsilon}(q_{\varepsilon},\dot{q}_{\varepsilon})$ can be written as

$$H_{\varepsilon}(q_{\varepsilon}, \dot{q}_{\varepsilon}) = \frac{1}{2} \dot{q}_{\varepsilon}^{T} M_{\varepsilon}(q_{\varepsilon}) \dot{q}_{\varepsilon} + V_{\varepsilon}(q_{\varepsilon}),$$

where $M_{\varepsilon}(q_{\varepsilon}) \in \mathbb{R}^{2n \times 2n}$ and $V_{\varepsilon}(q_{\varepsilon}) \in \mathbb{R}$ read as

$$M_{\varepsilon}(q_{\varepsilon}) = \begin{bmatrix} M_l(q_1) + I & -\varepsilon I \\ -\varepsilon I & \varepsilon^2 I \end{bmatrix}, \ V_{\varepsilon}(q_{\varepsilon}) = P_g(q_1) + \frac{1}{2}\varepsilon z^T z$$

Remark 2. Note that $H_0 = \lim_{\varepsilon \to 0} H_{\varepsilon}$ is precisely the Hamiltonian function of a rigid robot.

¹ Compactness ensures existence and uniqueness of S_{ε} , see Verhulst (2005).

(17)

Defining the generalized momenta as $p_{\varepsilon} = M_{\varepsilon}(q_{\varepsilon})\dot{q}_{\varepsilon}$, it is straightforward to rewrite the Hamiltonian as $H_{\varepsilon}(q_{\varepsilon}, p_{\varepsilon}) = \frac{1}{2}p_{\varepsilon}^{T}M_{\varepsilon}(q_{\varepsilon})^{-1}p_{\varepsilon} + V_{\varepsilon}(q_{\varepsilon})$.

Let us now perform the canonical transformation as briefly exposed in Section 2.2. For this, a change of coordinates $(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon}) = \Phi(q_{\varepsilon}, p_{\varepsilon})$ is defined by

$$\begin{bmatrix} \bar{q}_{\varepsilon} \\ \bar{p}_{\varepsilon} \end{bmatrix} = \begin{bmatrix} q_{\varepsilon} \\ \bar{T}_{\varepsilon} (\bar{q}_{\varepsilon})^{T} \dot{\bar{q}}_{\varepsilon} \end{bmatrix}, \tag{14}$$

where the matrix $\bar{T}_{\varepsilon}(\bar{q}_{\varepsilon}) \in \mathbb{R}^{2n \times 2n}$ is lower triangular and is defined by $M_{\varepsilon} = \bar{T}_{\varepsilon}(\bar{q}_{\varepsilon})\bar{T}_{\varepsilon}(\bar{q}_{\varepsilon})^T$. Under this change of coordinates the Hamiltonian is rewritten as $\bar{H}_{\varepsilon} = \frac{1}{2}\bar{p}_{\varepsilon}^T\bar{p}_{\varepsilon} + V(\bar{q}_{\varepsilon})$, and the port-Hamiltonian equations have the form (4), with $(M,\bar{T}) = (M_{\varepsilon},\bar{T}_{\varepsilon})$. In order to show that for a flexible-joint robot, the port-Hamiltonian equation takes the form of a slow-fast system we need to carefully study each term of (4). Let

$$\bar{T}_{\varepsilon}(\bar{q}_{\varepsilon}) = \begin{bmatrix} t_1 & 0 \\ t_2 & t_3 \end{bmatrix},$$

where $\bar{T}_{\varepsilon}(\bar{q}_{\varepsilon}) \in \mathbb{R}^{2n \times 2n}$, and for simplicity of notation we shall omit the dependence of t_i on \bar{q}_{ε} .

Remark 3. We assume that t_1 is known because of the relation $t_1t_1^T = M_l(q_1) + I$, which comes from a rigid robot model.

Notation. To save space, let us make the following definitions.

- First, let $\bar{q}_{\varepsilon} = (\bar{q}_1, \bar{q}_2) = (q_1, z) \in \mathbb{R}^{2n}$ and $\bar{p}_{\varepsilon} = (\bar{p}_1, \bar{p}_2) \in \mathbb{R}^{2n}$.
- $T_1 = T_1(\bar{q}) \in \mathbb{R}^{n \times n}$ is defined by $T_1 = t_1 t_1^T$.
- $t_4 = t_4(\bar{q}) \in \mathbb{R}^{n \times n}$ is defined by $t_4 t_4^T = I I(t_1 t_1^T)^{-1} I = I I T_1^{-1} I$, which we assume to be known since t_1 is known. It is a matter of simple linear algebra to show that t_4 exists and is unique.
- $\alpha = \alpha(\bar{q}) \in \mathbb{R}^{n \times n}$ is defined by $\alpha = t_4^{-1} I(t_1 t_1^T)^{-1}$.
- $\beta = \beta(\bar{q}, \bar{p}) \in \mathbb{R}^{n \times n}$ is defined by $\beta = \frac{\partial}{\partial \bar{q}_1} (t_1^{-1} \bar{p}_1)$.
- $\gamma = \gamma(\bar{q}, \bar{p}) \in \mathbb{R}^{n \times n}$ is defined by $\gamma = \frac{\partial}{\partial \bar{q}_1}(\alpha \bar{p}_1)$.

It follows from careful computations that

$$\bar{T}_{\varepsilon}^{-1} = \begin{bmatrix} t_1^{-1} & 0 \\ \alpha & \frac{1}{\varepsilon} t_4^{-1} \end{bmatrix}, \quad \bar{J}_{2,\varepsilon}(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon}) = \begin{bmatrix} j_1 & j_2 \\ -j_2^T & j_3 \end{bmatrix}, \quad (15)$$

where

$$j_1 = \beta t_1^{-T} - t_1^{-1} \beta^T$$

$$j_{2} = \underbrace{-\beta \alpha^{T} - t_{1}^{-1} \gamma^{T}}_{=:j_{21}} - \frac{1}{\varepsilon} \underbrace{t_{1}^{-1} \left(\frac{\partial}{\partial \bar{q}_{1}} (t_{4}^{-1} \bar{p}_{2}) \right)^{T}}_{=:j_{22}} = j_{21} - \frac{1}{\varepsilon} j_{22}$$

$$j_3 = \underbrace{-\gamma \alpha^T + \alpha \gamma^T}_{=:j_{31}} -$$

$$\frac{1}{\varepsilon} \left(\underbrace{\frac{\partial}{\partial \bar{q}_1} (t_4^{-1} \bar{p}_2) \alpha^T - \alpha \left(\frac{\partial}{\partial \bar{q}_1} (t_4^{-1} \bar{p}_2) \right)^T}_{=:j_{32}} \right)$$

$$=j_{31}-\frac{1}{\varepsilon}j_{32}$$

Remark 4. Note that $j_1 = -j_1^T$ and $j_3 = -j_3^T$ and therefore $\bar{J}_{2,\varepsilon}$ is indeed skew-symmetric.

Proposition 1. Under the coordinates $(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon})$ introduced above, a flexible-joint robot has the PH equations

$$\begin{bmatrix} \dot{\bar{q}}_1 \\ \dot{\bar{q}}_2 \\ \dot{\bar{p}}_1 \\ \dot{\bar{p}}_2 \end{bmatrix} = \begin{bmatrix} 0 & 0 & t_1^{-T} & \alpha^T \\ 0 & 0 & 0 & \frac{t_4^{-T}}{\varepsilon} \\ -t_1^{-1} & 0 & j_1 & j_{21} - \frac{j_{22}}{\varepsilon} \\ -\alpha^T & -\frac{t_4^{-1}}{\varepsilon} & -j_{21} + \frac{j_{22}}{\varepsilon} & j_{31} - \frac{j_{32}}{\varepsilon} \end{bmatrix} \begin{bmatrix} \frac{\partial \bar{H}_{\varepsilon}}{\partial \bar{q}_1} \\ \frac{\partial \bar{H}_{\varepsilon}}{\partial \bar{q}_2} \\ \frac{\partial \bar{H}_{\varepsilon}}{\partial \bar{p}_1} \\ \frac{\partial \bar{H}_{\varepsilon}}{\partial \bar{p}_2} \end{bmatrix} + \begin{bmatrix} 0_{n \times n} \\ 0_{n \times n} \\ \bar{G}_1 \\ \bar{G}_2 \end{bmatrix} \bar{v},$$

with $G_i = G_i(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon}, \varepsilon) \in \mathbb{R}^{n \times n}$, and where for notational simplicity we have left out the arguments of \bar{H}_{ε} , t_1 , t_4 , j_1 , j_{21} , j_{22} , j_{31} and j_{32} . Moreover, the corresponding reduced slow subsystem has the dynamics of a rigid-joint robot in the PH framework.

Proof. The PH equations are obtained by substituting (15) and (16) into (4). Regarding the reduced systems, note that we can rewrite the PH equations as

$$\dot{q}_{1} = t_{1}^{-} T \bar{p}_{1} + \alpha^{T} \bar{p}_{2}$$

$$\varepsilon \dot{q}_{2} = t_{4}^{-T} \bar{p}_{2}$$

$$\dot{\bar{p}}_{1} = -t_{1}^{-1} \frac{\partial \bar{H}_{\varepsilon}}{\partial \bar{q}_{1}} + j_{1} \bar{p}_{1} + \left(j_{21} - \frac{j_{22}}{\varepsilon}\right) \bar{p}_{2} + g_{1}(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon}) \bar{v}$$

$$\varepsilon \dot{\bar{p}}_{2} = -\varepsilon \alpha^{T} \frac{\partial \bar{H}}{\partial \bar{q}_{1}} - t_{4}^{-1} \frac{\partial \bar{H}}{\partial \bar{q}_{2}} - (\varepsilon j_{21} - j_{22}) \bar{p}_{1}$$

$$+ (\varepsilon j_{31} - j_{32}) \bar{p}_{2} + \varepsilon g_{2}(\bar{q}_{\varepsilon}, \bar{p}_{\varepsilon}) \bar{v}.$$
(18)

Taking now the limit $\varepsilon \to 0$ of (18), we find the constraints

$$0 = \bar{p}_2$$

$$0 = -t_4^{-1} \frac{\partial H_0}{\partial \bar{q}_2} + j_{22}\bar{p}_1 + j_{32}\bar{p}_2.$$

Furthermore, note that $\frac{\partial H_0}{\partial \bar{q}_2} = \varepsilon z$, $j_{22} \in O(\varepsilon)$, $\lim_{\varepsilon \to 0} q_{\varepsilon} = (q_1, 0)$, and $\lim_{\varepsilon \to 0} p_{\varepsilon} = ((M_l(q_1) + I)\dot{q}_1, 0)$. Therefore the reduced system is simply the rigid PH model

$$\begin{bmatrix} \dot{q}_1 \\ \dot{p}_1 \end{bmatrix} = \begin{bmatrix} 0 & t_1^{-T} \\ -t_1^{-1} & j_1 \end{bmatrix} \begin{bmatrix} \frac{\partial H_0}{\partial q_1} \\ \frac{\partial H_0}{\partial p_1} \end{bmatrix} + \begin{bmatrix} 0_{n \times n} \\ g_1(q_1, p_1) \end{bmatrix} u_s, \quad (19)$$

where u_s stands for the controller for the slow subsystem (12).

Remark 5. The PHS (19) is the PH model of a rigid robot, which physically means that the stiffness coefficient of the springs, of the flexible-joint robot, is infinity.

Since (17) has the structure of a slow-fast system ((with (\bar{q}_2, \bar{p}_2) being fast variables and (\bar{q}_1, \bar{p}_1) being the slow variables), we also find that the layer equation reads as

$$\begin{bmatrix} \bar{q}_2' \\ \bar{p}_2' \end{bmatrix} = \begin{bmatrix} 0 & t_4^{-T} \\ -t_4^{-1} & j_{32} \end{bmatrix} \begin{bmatrix} \frac{\partial \bar{H}}{\partial \bar{q}_2} \\ \frac{\partial \bar{H}}{\partial \bar{p}_2} \end{bmatrix} + \begin{bmatrix} 0_{n \times n} \\ g_2(\bar{q}_1, \bar{p}_1, \bar{q}_2, \bar{p}_2) \end{bmatrix} u_f,$$
(20)

where now (\bar{q}_1, \bar{p}_1) are fixed constants, and u_f stands for the controller of the layer equation (13). Note that the layer equation has also a PH structure.

With the previous exposition we have found reduced systems for a flexible-joint robot whose model is written in the PH framework. For control purposes, and following e.g., Kokotovic et al. (1986), it is possible to design controllers for a flexible-link robot from those for the rigid-robot and the layer equation.

4. NUMERICAL EXAMPLE

To illustrate our previous exposition, we present in this section a simulation of a composite control of a 2R planar flexible-joint manipulator. Joint flexibility can be attributed to several physical factors, like motor-to-link coupling, harmonic drives, etc Spong (1987); De Luca (2014). For simplicity, we assume that the robot acts on the horizontal plane and thus, gravitational effects are neglected. A schematic of the 2R flexible-joint robot is shown in Figure 1.

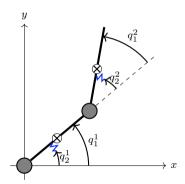


Fig. 1. Schematic of a 2R planar, flexible-joint robot.

Let $q_1=(q_1^1,q_1^2)$ and $q_2=(q_2^1,q_2^2)$ be the coordinates of the links and of the motors, respectively. Each link has length l_i , mass m_i , inertia I_{l_i} and distance to the center of mass r_i ; while each motor has associated inertia I_i , i=1,2. We assume that the matrix $I\in\mathbb{R}^{2\times 2}$, such that $I=\operatorname{diag}\{I_1,I_2\}$, and the matrix $K\in\mathbb{R}^{2\times 2}$ is also diagonal of the form $K=\operatorname{diag}\{I_{l_i},I_{l_i}\}$, thus fitting in the exposition above. Let us define the following constants

$$a_1 = m_1 r_1^2 + m_2 l_1^2 + I_{l_1}$$

 $a_2 = m_2 r_2^2 + I_{l_2}$
 $b = m_2 l_1 r_2$.

Then, the matrix $M_l(q_1) \in \mathbb{R}^{2 \times 2}$ reads as

$$M_l(q_1) = \begin{bmatrix} a_1 + a_2 + 2b\cos q_1^2 & a_2 + b\cos q_1^2 \\ a_2 + b\cos q_1^2 & a_2 \end{bmatrix}.$$

The task is to make the *links* (q_1^1, q_1^2) follow a desired trajectory $\bar{q}_{1,d} = q_{1,d} = (q_{1,d}^1, q_{1,d}^2)$ given by $q_{1,d}^1 = q_{1,d}^2 = 0.1 + 0.05 \sin(t)$. To achieve such a task, we implement separate controllers following the design principle of Kokotovic

et al. (1986). This is: one controller is designed for the rigid robot (the slow subsystem (19)) independently of the fast subsystem, and another controller is designed for the fast subsystem, where now the slow variables are taken as fixed parameters. These two controllers shall guarantee stability in their own domain (slow and fast reduced subsystems respectively). Then, the controller for the flexible-joint robot is defined as the sum of both reduced controllers. The stability of the flexible system is guaranteed by GSPT arguments, recall Section 2.3 and see Kokotovic (1984); Fenichel (1979); Jones (1995); Kaper (1999).

The controller synthesis for the rigid robot, which has a PH equation of the form (19) is taken from Dirksz and Scherpen (2013) and reads as

$$u_{s} = M_{l}\ddot{q}_{1,d} + \frac{\partial}{\partial q_{1}} (M_{l}\dot{q}_{1,d})\dot{q}_{1,d} - \frac{1}{2} \frac{\partial}{\partial q_{1}} (\dot{q}_{1,d}^{T} M_{l}\dot{q}_{1,d}) - K_{p}(q_{1} - q_{1,d}) - K_{c}(q_{1} - q_{1,d} - q_{1,c}),$$
(21)

where q_c is the controller state and its dynamics are given by

$$\dot{q}_{1,c} = K_d^{-1} K_c (q_1 - q_{1,d} - q_{1,c}),$$

and K_c , K_d and K_p are positive definite matrices, see Dirksz and Scherpen (2013) for more details.

Remark 6. Even though (19) is not exactly the same model as considered in Dirksz and Scherpen (2013), both are, up to the change of coordinates of Section 2.2, a standard mechanical system. Moreover, note that in the change of coordinates (14) $\bar{q}_{\varepsilon}=q_{\varepsilon}.$

The controller u_s applied to the rigid robot has the performance shown in Figure 2a.

For the fast subsystem, which has the form (20), we employ the same controller design idea of Dirksz and Scherpen (2013). This is due to the fact that the layer subsystem is also a PH system, just now we have a desired trajectory $\bar{q}_{2,d} = z_d = (z_d^1, z_d^2) = (0,0)$. The reason is that we want to follow the desired trajectory with a zero deflection, i.e. $z = q_1 - q_2 = 0$. This yields a controller of the form

$$u_f = -L_p z - L_c (z - z_c),$$
 (22)

with the controller dynamics

$$\dot{z}_c = L_d^{-1} L_c(z - z_c),$$

where L_d , L_p and L_c are positive definite matrices.

Remark 7. Note that the controllers u_s and u_f of (21) and (22), respectively, only use position measurements.

By combining these two controllers as $u=u_s+u_f$, and implementing them into the flexible-joint robot whose model is now of the form (17), we get the performance shown in Figures 2b and 2c. From these we see that the robot closely follows the desired trajectories after one second. Finally, comparing Figures 2a and 2b we note that the difference between the rigid and the flexible robot behaviors is barely noticeable.

5. CONCLUSIONS

In this document we have explored the methodology of model order reduction based on singular perturbations for a PH system. In order to do so, we have written the PH model of a flexible-joint robot in a slow-fast format.

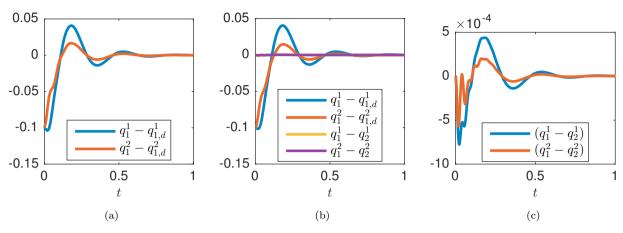


Fig. 2. (a) The trajectory tracking error for a rigid 2R robot. (b) The trajectory tracking error for a 2R flexible-joint robot. For this simulation $\varepsilon = 0.01$ or $K = \text{diag}\{100, 100\}$. (c) Zoomed-in deflection.

Then, by inspecting the structure of the system we have shown, as it is to be expected, that the corresponding slow subsystem is the PH model of a rigid-joint robot. Consequently, as it happens in the EL framework, the design of controllers from the reduced subsystems is justified. To exemplify our exposition, we have implemented a controller with only position measurements designed in Dirksz and Scherpen (2013). Our simulations show a good performance of the controllers. However, we have used controllers designed only on the reduced subsystems. The performance of the closed-loop systems can be improved by taking higher order terms of ε when designing the controller. In particular, and as a natural extension of this work, a careful study of the fast subsystem is required in order to rigorously prove exponential stability under composite controllers within the PH framework. Moreover, as future research, visco-elastic joints and elastic links may be incorporated to the PH model.

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